



## RESEARCH ARTICLE OPEN ACCESS

Triplet 4-Nitrenemethyl-Pyridine-*N*-Oxide: A Model Alkyl Nitrene Isolated in Cryogenic MatricesCláudio M. Nunes<sup>1</sup>  | A. J. Lopes Jesus<sup>2</sup>  | Tatiana Caneca<sup>1</sup>  | Gemma Aragay<sup>3</sup>  | Pablo Ballester<sup>3,4</sup>  | Rui Fausto<sup>1,5</sup> <sup>1</sup>University of Coimbra, CQC-IMS, Department of Chemistry, Coimbra, Portugal | <sup>2</sup>CQC-IMS, Faculty of Pharmacy, Coimbra, Portugal | <sup>3</sup>Institute of Chemical Research of Catalonia (ICIQ), the Barcelona Institute of Science and Technology (BIST), Tarragona, Spain | <sup>4</sup>ICREA, Barcelona, Spain | <sup>5</sup>Spectroscopy@IKU, Faculty Sciences and Letters, Department of Physics, Istanbul Kultur University, Istanbul, Turkey**Correspondence:** Cláudio M. Nunes ([cmnunes@qui.uc.pt](mailto:cmnunes@qui.uc.pt))**Received:** 13 January 2026 | **Revised:** 5 February 2026 | **Accepted:** 7 February 2026**Keywords:** IR spectroscopy | matrix-isolation | nitrenes | photochemistry | quantum chemical calculations**ABSTRACT**

Alkyl nitrenes are highly reactive intermediates that are particularly challenging to study experimentally. Here, we report the first investigation of the simplest triplet alkyl nitrene bearing a pyridine-*N*-oxide moiety. Direct irradiation ( $\lambda > 350$  nm) of 4-azidomethyl-pyridine-*N*-oxide in an Ar matrix (15 K) generated the target triplet 4-nitrenemethyl-pyridine-*N*-oxide **3**<sup>2</sup>, along with *E* and *Z* isomers of 4-iminomethyl-pyridine-*N*-oxide **3**. Results from subsequent irradiations ( $\lambda > 325$  nm) enabled the discrimination of the IR spectroscopic signatures of these species, allowing their unequivocal identification with support from DFT-B3LYP computations. Triplet nitrene **3**<sup>2</sup> was found to be stable in dark Ar matrices, demonstrating that quantum tunneling 1,2-H shift to imine **3** is not operative on the experimental timescale. Computations show that this reaction is thermodynamically favorable on the triplet surface but involves a high-energy barrier ( $\sim 36$  kcal mol<sup>-1</sup>), explaining the absence of H-tunneling. As the pyridine-*N*-oxide moiety can act as an anchoring site for complexation with calix[4]pyrrole derivatives, our findings open new perspectives for investigating alkyl nitrene reactivity and tunneling phenomena under supramolecular confinement conditions.

**1 | Introduction**

Nitrenes are reactive intermediates characterized by a monovalent nitrogen center bearing six valence electrons, whose properties strongly depend on their molecular structure and spin multiplicity [1, 2]. Among the various classes of nitrenes, alkyl nitrenes remain particularly challenging to study experimentally. In contrast to their aryl counterparts, the lack of significant electronic delocalization at the nitrene center renders alkyl nitrenes extremely reactive and difficult to characterize.

Attempts to generate the simple methylnitrene [CH<sub>3</sub>N] via thermal or photochemical denitrogenation of methyl azide [CH<sub>3</sub>N<sub>3</sub>] have eluded both time-resolved spectroscopic detection and trapping by bimolecular reagents [1]. Direct photolysis

experiments of methyl azide in cryogenic matrices could only identify the corresponding methyleneimine [CH<sub>2</sub>=NH]<sup>1,2</sup> [3–7]. Methylnitrene was characterized as possessing a triplet ground state that lies  $\sim 31.2$  kcal mol<sup>-1</sup> lower in energy than the singlet state [8]. Computations at the CASPT2 level indicate that singlet methylnitrene is a true minimum on the potential energy surface (PES), although it is separated by a very low barrier of 2–3 kcal mol<sup>-1</sup> from the 1,2-hydrogen shift to methyleneimine, with an exothermicity exceeding 80 kcal mol<sup>-1</sup> [9]. Supporting this rationalization, a collision-free photochemistry study of methyl azide (248 nm), probed by tunable synchrotron radiation-based photoionization, indicated direct N<sub>2</sub> loss and singlet methylnitrene formation, which rapidly rearranges to yield the vibrationally excited methanimine precluding intersystem crossing (ISC) to the triplet state [10].

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## 2.2 | Matrix-Isolation IR Spectroscopy

A sample of azide **1** was placed in a custom-built oven connected to the vacuum chamber of a closed-cycle helium cryostat (APD Cryogenics model HC-2, equipped with a DE-202A expander) and pumped to remove any volatile impurities. Prior to matrix deposition, the chamber pressure reached  $\sim 2 \times 10^{-6}$  mbar. Vapors of **1** were generated by gently heating the sample to  $\sim 45^\circ\text{C}$  using a thermoelectric power supply (Vitecom DC Power Supply, model 75-HY5003). The resulting vapors were co-deposited with a large excess of high-purity argon (N60, Air Liquide) onto a CsI optical window maintained at 15 K. The temperature was measured directly at the sample holder by using a silicon diode sensor connected to a digital controller (Scientific Instruments, model 9650-1), providing a temperature stabilization with an accuracy of  $\pm 0.1$  K. Matrix preparation was monitored by recording IR spectra to optimize deposition and ensure conditions for isolating a monomeric sample. IR spectra were collected in the mid-infrared region ( $4000\text{--}400\text{ cm}^{-1}$ ) at a resolution of  $0.5\text{ cm}^{-1}$  using a Thermo Nicolet 6700 FTIR spectrometer equipped with a deuterated triglycine sulfate (DTGS) detector and a germanium/potassium bromide (Ge/KBr) beamsplitter. The instrument was continuously purged by a stream of dry air to minimize interference from moisture and  $\text{CO}_2$ .

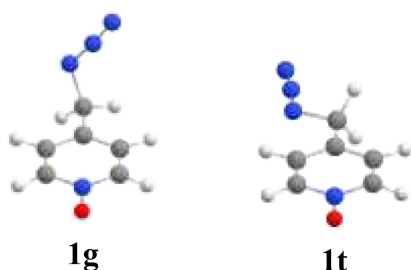
## 2.3 | UV-Vis Irradiation Experiments

UV-Vis irradiation was conducted through the outer quartz window of the cryostat using a 250 W Hg (Xe) arc lamp (Newport, Oriel Instruments). The lamp output was passed through an 8 cm water filter and long-pass optical filters with cutoff wavelengths of 350 or 325 nm.

**TABLE 1** | Optimized geometries of conformers **1g** and **1t**, relative zero-point-corrected energies ( $\Delta H_{0\text{K}}$ ,  $\text{kJ mol}^{-1}$ ), relative Gibbs free energies at 298 K ( $\Delta G_{298\text{K}}$ ,  $\text{kJ mol}^{-1}$ ), and the respective Boltzmann populations at the same temperature (*Pop.*, %), computed at the B3LYP/6-311+G(2d,p) and DLPNO-CCSD(T)/def2-TZVPP// $\omega$ B97XD/6-311+G(2d,p) levels, respectively<sup>a</sup>.

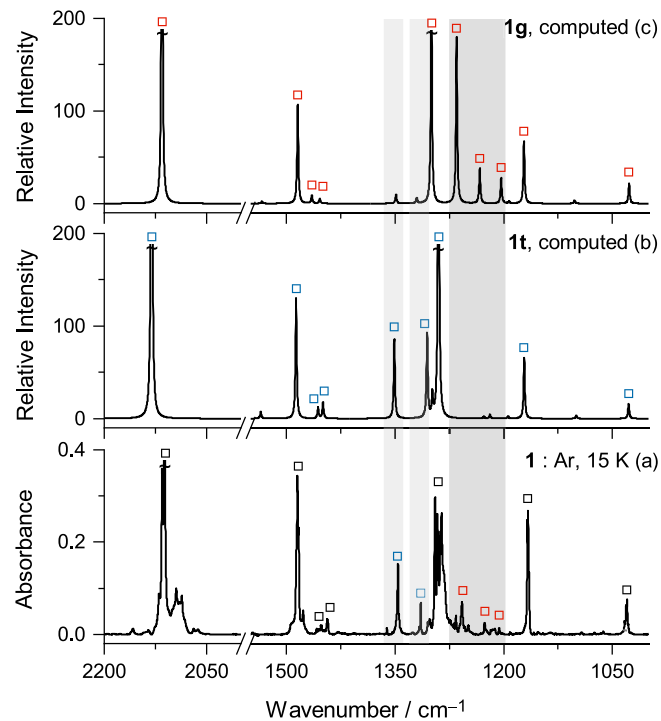
	<b>1g</b>	<b>1t</b>
$\Delta H_{0\text{K}}$	0.00   0.00	0.41   1.94
$\Delta G_{298\text{K}}$	0.00   0.00	0.29   0.55
<i>Pop.</i>	54   56	46   44

<sup>a</sup>Relative Gibbs free energies were corrected for conformational degeneracy by including the term  $-RT \ln(g)$ , where  $g$  is the number of equivalent minima associated with each conformer ( $g=4$  for **1g** and  $g=2$  for **1t**; see Figure S1). Color scheme: carbon (gray), nitrogen (blue), oxygen (red), and hydrogen (white).



## 2.4 | Computations

A two-dimensional relaxed PES scan was performed for azide **1** using the DFT-B3LYP [29–31] functional in conjunction with the 6-311+G(2d,p) basis set [32]. The procedure involved stepwise variation of the two dihedral angles defining the orientation of the azidomethyl group, although all other internal coordinates were optimized (Figure S1). The minimum energy structures identified for **1** and for the photoproducts investigated in this work were fully optimized and subsequently subjected to harmonic vibrational frequency calculations at the same level of theory. The resulting vibrational wavenumbers were scaled by factors of 0.979 below  $3000\text{ cm}^{-1}$  and 0.960 above  $3000\text{ cm}^{-1}$  and, together with the corresponding IR intensities, were used for spectral simulations [33]. Each band was convoluted with a Lorentzian function with a full width at half maximum (FWHM) of  $2\text{ cm}^{-1}$ , centered at the scaled wavenumbers, using the ChemCraft software (Version 1.6) [34]. The integral band intensities correspond to the calculated absolute IR intensities and are reported in arbitrary units as relative intensities. Vibrational mode assignments were performed using vibAnalysis software [35] and by visual inspection of normal-mode animations. To obtain more accurate energies, the relevant species and transition states were subjected to geometry optimizations and harmonic vibrational frequency calculations at the  $\omega$ B97XD/6-311+G(2d,p) level [36], and single-point energies were subsequently computed at the DLPNO-CCSD(T)/def2-TZVPP level [37–40].  $T_1$  diagnostics from DLPNO-CCSD(T) calculations

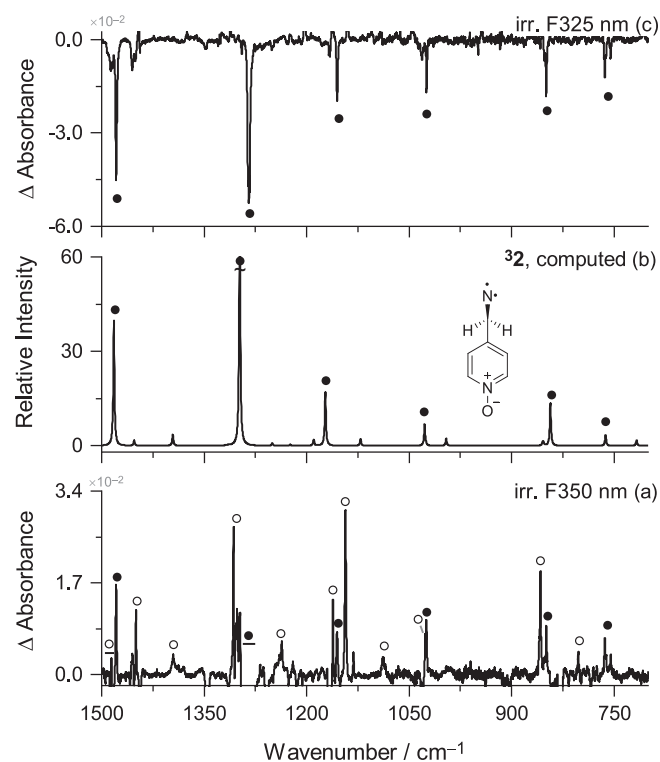


**FIGURE 1** | (a) Experimental IR spectrum of 4-azidomethyl-pyridine-*N*-oxide **1** isolated in an Ar matrix at 15 K. (b,c) B3LYP/6-311+G(2d,p) computed IR spectra of *trans* and *gauche* conformers of azide **1** (**1t** and **1g**). The light gray and dark gray backgrounds indicate bands assigned to conformers **1t** (blue squares) and **1g** (red squares), respectively.

(0.016 for **32** and 0.023 for **33** and the corresponding transition state) indicate only limited multireference character, consistent with a reliable coupled-cluster description of these open-shell systems. All DFT computations were carried out with Gaussian 16 (Revision B.01) [41], whereas the coupled-cluster single-point calculations were performed using the ORCA (Version 6.1.0) [42–45]. Hirshfeld atomic spin populations of triplet nitrene **32** and triplet imine **33E** were calculated from B3LYP/6-311+G(2d,p) spin densities using the Multiwfn program [46]. Spin contamination in the B3LYP descriptions of the open-shell triplet species was found to be negligible based on  $\langle S^2 \rangle$  values close to the ideal value of 2.000.

### 3 | Results and Discussion

The conformational space of 4-azidomethyl-pyridine-*N*-oxide **1** was explored at the B3LYP/6-311+G(2d,p) level as a function of rotation around the exocyclic C–C and C–N bonds (Figure S1). Two distinct minima were identified, conformers **1g** (fourfold degenerate) and **1t** (doubly degenerate) (Table 1). In **1g**, the C–N bond is tilted out of the ring plane, and the azide group adopts *gauche* (**g**) orientation relative to the C–C bond. In **1t**, the C–N bond lies approximately in the ring plane, with the azide group



**FIGURE 2** | (a) Experimental difference IR spectrum showing changes after irradiation at  $\lambda \geq 350$  nm (250 W, 2 min) of 4-azidomethyl-pyridine-*N*-oxide **1** in an Ar matrix at 15 K. Negative bands due to consumption of **1** are truncated. (b) B3LYP/6-311+G(2d,p) computed IR spectrum of triplet 4-nitrenemethyl-pyridine-*N*-oxide **32**. (c) Experimental difference IR spectrum showing changes after irradiation at  $\lambda \geq 325$  nm (250 W, 12 min), subsequent to the consumption of **1** at  $\lambda \geq 350$  nm (250 W, 1 h). Positive bands due to unidentified products are truncated. Closed circles indicate bands of the photoproduct **A** assigned to **32**, and open circles indicate bands of the photoproduct **B**.

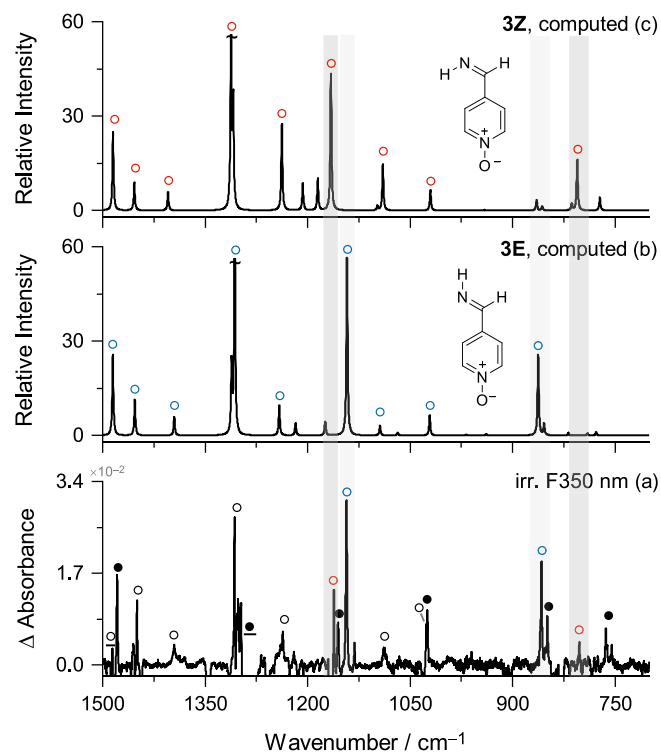
**TABLE 2** | Experimental IR spectral data (Ar matrix 15 K), B3LYP/6-311+G(2d,p) computed wavenumbers ( $\tilde{\nu}/\text{cm}^{-1}$ ) and IR intensities ( $\text{I}/\text{kmol}^{-1}$ ), and approximate vibrational assignment of triplet 4-nitrenemethyl-pyridine-*N*-oxide **32**.

Experimental <sup>a</sup>	Calculated <sup>b</sup>		Approximate assignment <sup>c</sup>
Ar, 15 K	$\tilde{\nu}$	I	
1624 (w)*	1626	36.6	$\nu$ CC ring
1479 (s)	1482	125.1	$\delta$ CH ring; $\nu$ NO – $\nu$ CC exo
n.o.	1396	11.3	$\delta$ CH <sub>2</sub>
1286/1283 (s)	1298	252.7	$\nu$ NO
1156 (m)	1173	53.6	$\nu$ CC ring; $\delta$ CH ring
1025 (m)	1027	21.4	$\nu$ CN ring; $\delta$ ring
852/849 (m)	843	42.4	$\gamma$ CH ring; $\gamma$ ring
763/755 (w)	762	10.6	$\nu$ CC exo; $\delta$ ring

<sup>a</sup>Abbreviation “n.o.” refers to a non-observed band. Asterisk symbol refers to a band partially overlapped with monomeric water. Experimental intensities are expressed in qualitative terms: s, strong; m, medium; w, weak.

<sup>b</sup>Only vibrational modes in the 1800–550  $\text{cm}^{-1}$  region having calculated intensities  $\geq 10 \text{ km mol}^{-1}$  were considered. Harmonic wavenumbers were scaled by 0.979.

<sup>c</sup>Approximate assignment was based on the results provided by the “vibAnalysis” software and supported by ChemCraft animation of the corresponding vibrations. Abbreviations:  $\nu$  = stretching,  $\delta$  = in-plane bending,  $\gamma$  = out-of-plane bending; exo = exocyclic fragment. The symbol “–” denotes vibrational combinations occurring in the opposite phase (out-of-phase).



**FIGURE 3** | (a) Experimental difference IR spectrum showing changes after irradiation at  $\lambda \geq 350$  nm (250 W, 2 min) of 4-azidomethyl-pyridine-*N*-oxide **1** in an Ar matrix at 15 K. Negative bands due to consumption of **1** are truncated. (b,c) B3LYP/6-311+G(2d,p) computed IR spectra of *E* and *Z* isomers of 4-iminomethyl-pyridine-*N*-oxide (**3E** and **3Z**). The light gray and dark gray backgrounds indicate bands assigned specifically to isomer **3E** (blue open circles) and **3Z** (red open circles), respectively.

assuming *trans* (**t**) configuration with respect to the C–C bond. The two conformers are nearly isoenergetic, which is confirmed at the higher DPLNO-CCSD(T)/def2-TZVPP// $\omega$ B97XD/6-311+G(2d,p) level. Boltzmann populations calculated from the Gibbs free energies of **1** computed in the gas phase at room temperature indicate that conformers **1g** and **1t** are present in similar proportions (Table 1).

The synthesized 4-azidomethyl-pyridine-*N*-oxide **1** precursor of **32** was gently evaporated slightly above room temperature and its vapors were co-deposited with a large excess of Ar onto the cryostat optical window cooled to 15 K. The resulting IR spectrum of monomers of **1** isolated in an Ar matrix reveals the presence of both **1g** and **1t** conformers (Figure 1). In particular, the

IR bands at 1346 and 1315  $\text{cm}^{-1}$  are in close agreement with the B3LYP computed vibrational modes of **1t** at 1349 and 1306  $\text{cm}^{-1}$ , whereas the IR bands at 1258, 1227, and 1206  $\text{cm}^{-1}$  are exclusively compatible with the computed vibrational modes of **1g** at 1265, 1233 and 1204  $\text{cm}^{-1}$ . A comprehensive assignment of the IR spectrum of **1** is given in Table S1.

The spectroscopic data clearly indicate that in the deposited matrix, the population of **1t** is significantly larger than that of **1g**. These results can be interpreted as follows: Although the computed zero-point-corrected energies suggest that **1g** is marginally more stable than **1t** in the gas phase, under matrix-isolation conditions, the nearly planar **1t** becomes the most energetically favored form due to its better accommodation within the packed

**TABLE 3** | Experimental IR spectral data (Ar matrix 15 K), B3LYP/6-311+G(2d,p) computed wavenumbers ( $\tilde{\nu}/\text{cm}^{-1}$ ) and IR intensities ( $I/\text{km mol}^{-1}$ ), and approximate vibrational assignment of *E* and *Z* isomers of 4-iminomethyl-pyridine-*N*-oxide, **3E** and **3Z**.

Experimental <sup>a</sup>	Calculated <sup>b</sup>		Isomer	Approximate assignment <sup>c</sup>
	Ar, 15 K	$\tilde{\nu}$		
1639 (w)	1646/1644	65.9/38.5	<b>3E/3Z</b>	$\nu\text{C=N}$
1614 (m)	1623/1620	177.7/183.3	<b>3E/3Z</b>	$\nu\text{CC ring}; \nu\text{C=N}$
1541 (w)	1529/1528	19.0/12.9	<b>3E/3Z</b>	$[\nu\text{CC}; \nu\text{CN}]$ ring
1486 (ov)	1485/1485	80.9/78.0	<b>3E/3Z</b>	$\delta\text{CH ring}; \nu\text{NO} - \nu\text{CC exo}$
1450 (m)	1454/1453	27.8/35.7	<b>3Z/3E</b>	$[\nu\text{CC}; \delta\text{CH}]$ ring
1395 (br)	1404/1395	18.6/18.5	<b>3Z/3E</b>	$[\delta\text{NH}; \delta\text{CH}]$ exo
1307 (s)	1312/1311	189.3/69.4	<b>3Z/3E</b>	$\nu\text{NO}/\delta\text{CH ring}$
1303/1302/1297 (m)	1309/1307	104.0/210.9	<b>3Z/3E</b>	$\delta\text{CH ring}/\nu\text{NO}$
1236 (m)	1241/1238	30.2/86.7	<b>3E/3Z</b>	$\nu\text{CN ring}; [\delta\text{NH}; \delta\text{CH}]$ exo; $\nu\text{CC exo}$
1220 (w)	1218/1207	12.5/27.3	<b>3E/3Z</b>	$[\nu\text{CC}; \nu\text{CN}]$ ring
n.o.	1185	31.4	<b>3Z</b>	$\delta\text{NH}; \nu\text{CC exo}$
n.o.	1174	13.9	<b>3E</b>	$\delta\text{NH}; \delta\text{CH ring}$
1161 (m)	1166	136.3	<b>3Z</b>	$[\nu\text{CC}; \delta\text{CH}]$ ring
1143 (s)	1142	178.6	<b>3E</b>	$\delta\text{NH}; \nu\text{CC exo}; \delta\text{CH ring}$
1088 (m)	1094/1090	10.0/46.0	<b>3E/3Z</b>	$\delta\text{CH ring}/[\gamma\text{NH} + \gamma\text{CH}]$ exo
1026 (m)	1021/1020	20.3/20.2	<b>3E/3Z</b>	$\nu\text{CN ring}; \delta$ ring
n.o.	865	10.7	<b>3Z</b>	$\gamma$ ring; $\gamma\text{C}_{\text{exo}}; \gamma\text{CH ring}; \gamma\text{CH exo}; \gamma\text{NH}$
858 (s)	862	80.5	<b>3E</b>	$\gamma$ ring; $\gamma\text{CH ring}; \gamma\text{NH}$
	854	11.3	<b>3E</b>	$\gamma$ Ring breathing
802 (w)	805	51.0	<b>3Z</b>	$\gamma$ ring; $\gamma\text{C}_{\text{exo}}; \gamma\text{CH ring}; \gamma\text{CH exo}; \gamma\text{NH}$
n.o.	772	13.2	<b>3Z</b>	$\nu\text{CC exo}; \delta$ ring
n.o.	691	19.9	<b>3Z</b>	$\gamma$ ring; $\gamma\text{C}_{\text{exo}}; \gamma\text{CH ring}; \gamma\text{CH exo}; \gamma\text{NH}$
593 (m)	596/586	63.1/21.4	<b>3E/3Z</b>	$\delta\text{CC=N}; \delta$ ring

<sup>a</sup>Abbreviation "n.o." refers to a non-observed band. Experimental intensities are expressed in qualitative terms: s, strong; m, medium; br = broad; w, weak; ov = overlapped.

<sup>b</sup>Only vibrational modes in the 1800–550  $\text{cm}^{-1}$  region having calculated intensities  $\geq 10 \text{ km mol}^{-1}$  were considered. Harmonic wavenumbers were scaled by 0.979.

<sup>c</sup>Approximate assignment was based on the results provided by the "vibAnalysis" software and supported by ChemCraft animation of the corresponding vibrations. Abbreviations:  $\nu$  = stretching,  $\delta$  = in-plane bending,  $\gamma$  = out-of-plane bending; exo = exocyclic fragment. The symbols "+" and "-" denote vibrational combinations occurring in the same phase (in-phase) and in opposite phases (out-of-phase), respectively.

matrix lattice [47–49]. Given the very low energy barrier separating these conformers ( $\sim 3\text{--}5\text{ kJ mol}^{-1}$ ; see Figure S2), **1g** partially converts to the lower-energy **1t** during the matrix deposition. This phenomenon, often referred to as conformational cooling, leads to a matrix population of conformational isomers that differs from the original gas-phase equilibrium, which would otherwise be “frozen” under matrix-isolation conditions [50–52].

Irradiation of matrix-isolated azide **1**, using a 350-nm long-pass cutoff filter, results in two primary photoproducts, **A** and **B**, which are the only species detected within the first 2 min ( $\sim 17\%$  consumption of **1**) (Figure 2a). Upon prolonged irradiation under these conditions, one of the primary products (**B**) decreases its amount in the matrix concomitantly with the continued consumption of **1**, whereas the other (**A**) keeps accumulating (Figures S3 and S4). The secondary photoproducts arising from the consumption of **B** remain unidentified. Subsequent irradiation at shorter wavelengths, using a 325-nm long-pass cutoff filter, consumes the accumulated primary product **A**, thereby enabling a relatively cleaner IR spectral signature to be recorded. The good agreement between the experimental and B3LYP-computed spectra allows us to identify **A** as triplet 4-nitrenemethyl-pyridine-*N*-oxide **32** (Figure 2c,b). For instance, the most intense IR bands of **32** are observed at 1479 and  $\sim 1285\text{ cm}^{-1}$  and are well reproduced by the most intense computed IR absorption at 1482 and  $1298\text{ cm}^{-1}$ , respectively. Except for the computed vibrational mode at  $1396\text{ cm}^{-1}$  ( $I = 11.3\text{ km mol}^{-1}$ ), all other modes of **32** with appreciable intensity ( $\geq 10\text{ km mol}^{-1}$ ) were successfully identified in the experimental spectrum, as detailed in Table 2.

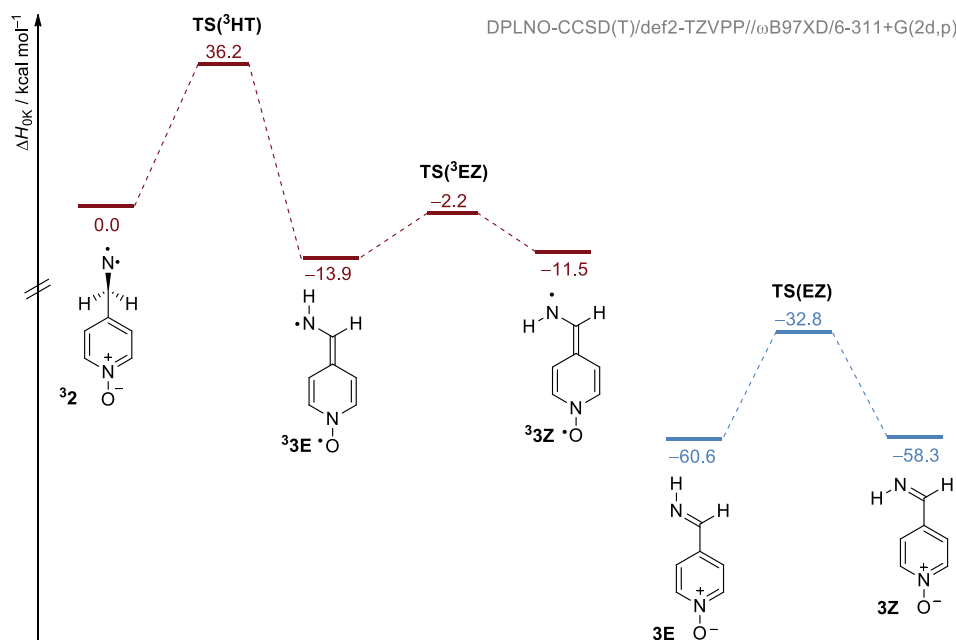
Having unambiguously assigned primary product **A** as the triplet nitrene **32**, the remaining bands in the IR spectrum recorded after the first 2 min of irradiation of **1** correspond to the spectral signature of the primary product **B** (Figures 2a and 3a). On the

basis of these experimental data and supported by B3LYP computations, we identify **B** as 4-iminomethyl-pyridine-*N*-oxide **3**, which exists in two isomeric forms, *E* (**3E**) and *Z* (**3Z**) (Figure 3). For instance, the experimental IR bands observed at 1143 and  $858\text{ cm}^{-1}$  show excellent correspondence with the computed IR absorption of **3E** at 1142 and  $862/854\text{ cm}^{-1}$ , respectively. In contrast, the experimental IR bands observed at 1161 and  $802\text{ cm}^{-1}$  agree well with the computed IR modes of **3Z** at 1166 and  $805\text{ cm}^{-1}$ , respectively. For most other regions, the experimental IR spectrum agrees well with the B3LYP-computed spectra of both **3E** and **3Z**, which differ only slightly. A detailed assignment of the IR spectra identifying the two imine isomers **3E** and **3Z** is given in Table 3.

In an additional experiment, after photogeneration of triplet nitrene **32** from **1** in an Ar matrix, the sample was kept in the dark for 24 h. IR monitoring revealed no detectable changes, demonstrating that **32** is stable under cryogenic matrix conditions over this period. This observation strongly suggests that quantum tunneling 1,2-H shift from **32** to imine **3** is not operative on the experimental timescale.

To rationalize these results, we explored the PES connecting triplet nitrene **32** with the corresponding triplet imine **33** (Figure 4). DPLNO-CCSD(T)// $\omega$ B97XD calculations reveal that **33** is thermodynamically more stable than **32** by more than  $10\text{ kcal mol}^{-1}$ . Consequently, the persistence of triplet nitrene **32** in the matrix and the absence of a tunneling-driven 1,2-H shift to **33** cannot be rationalized on thermodynamic grounds, given the exothermic nature of the rearrangement.

Spin density analysis provides insights into the origin of the relative energy of **32** and **33** species. As expected for an alkyl nitrene, the spin density in **32** is highly localized at the nitrene nitrogen atom, with only minimal delocalization onto the molecular



**FIGURE 4** | Reaction energy profile ( $\Delta H_{0K}$  in  $\text{kcal mol}^{-1}$ ) computed at the DPLNO-CCSD(T)/def2-TZVPP// $\omega$ B97XD/6-311+G(2d,p)+ZPVE level of theory. Values are relative to the energy of triplet nitrene **32**.

framework (Figure 5 and Table 4). In contrast, the triplet imine  $^3\mathbf{3}$  exhibits pronounced spin delocalization over the exocyclic nitrogen, the *ipso* carbon, and the *N*-oxide fragment. This extensive delocalization, strongly enhanced by the *N*-oxide moiety, provides a compelling explanation for the surprising relative stabilization of triplet imine  $^3\mathbf{3}$  (see comparison with the energetics of the benzyl and ethyl derivatives in Figure S5). Consistent with this picture, the computed energy barrier for *E/Z* isomerization of  $^3\mathbf{3}$  is only  $\sim 10 \text{ kcal mol}^{-1}$ , indicating that the exocyclic C–N bond possesses significant single-bond character rather than that of a typical imine double bond. Altogether, these data indicate that triplet imine  $^3\mathbf{3}$  is better described as a delocalized diradicaloid species.

Although the rearrangement from  $^3\mathbf{2}$  to  $^3\mathbf{3}$  is thermodynamically allowed, computations predict that this 1,2-H shift has a high-energy barrier of  $\sim 36 \text{ kcal mol}^{-1}$  (Figure 4). Whereas a few extraordinary examples of H-tunneling reactions through significant energy barriers ( $\sim 20\text{--}30 \text{ kcal mol}^{-1}$  range) [18, 53–56] have been reported in noble-gas cryogenic matrices, no tunneling reactions have been observed for barriers on the order of  $\sim 36 \text{ kcal mol}^{-1}$ . Such a high barrier is therefore expected to effectively suppress 1,2-H shift tunneling from  $^3\mathbf{2}$  to  $^3\mathbf{3}$  under cryogenic conditions<sup>3</sup>. At higher temperatures, however, thermally assisted tunneling may become feasible [53, 57], an intriguing possibility that merits further investigation.

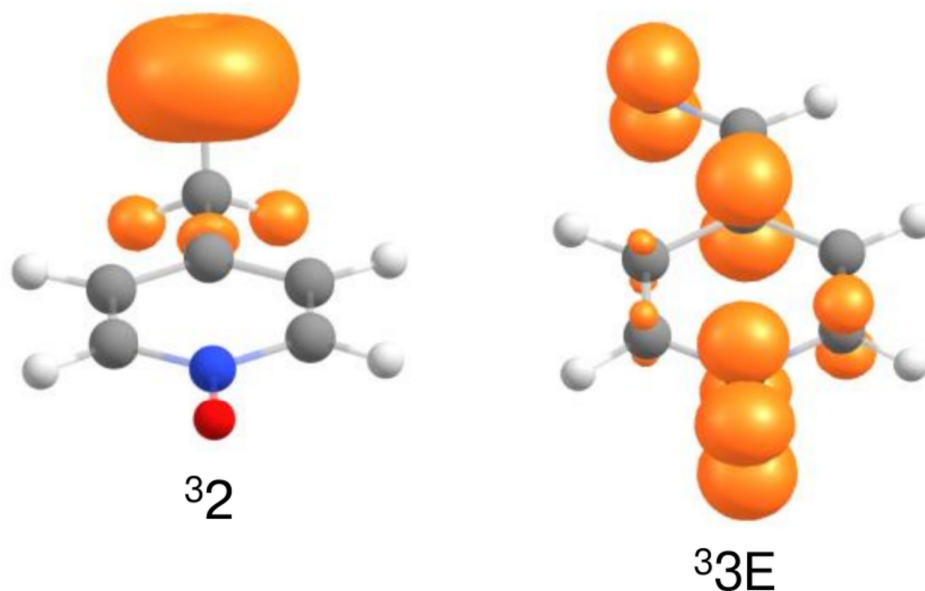
Finally, we can exclude the formation of imine  $\mathbf{3}$  as a secondary product arising from the photochemistry of nitrene  $^3\mathbf{2}$ . Monitoring the photolysis of  $\mathbf{1}$  as a function of time showed that  $\mathbf{3}$  reaches its maximum amount within the first minutes, whereas  $^3\mathbf{2}$  continues to accumulate in the matrix (Figure S4). If imine  $\mathbf{3}$  were formed via nitrene  $^3\mathbf{2}$ , its maximum amount in the matrix would be expected to follow that of  $^3\mathbf{2}$ . Therefore,  $\mathbf{3}$  is most likely formed in competition with ISC to  $^3\mathbf{2}$  from

photolysis of azide  $\mathbf{1}$  (Scheme 2). Assuming a singlet–triplet energy gap for nitrene  $\mathbf{2}$  comparable to that of methylnitrene ( $\sim 30 \text{ kcal mol}^{-1}$ ) [8], the corresponding 1,2-H shift from  $^1\mathbf{2}$  to imine  $\mathbf{3}$  would be highly exothermic, exceeding  $80 \text{ kcal mol}^{-1}$ , as reported for methylnitrene [9]. The reason why ISC can effectively compete with the ultrafast 1,2-H shift on the singlet surface, enabling the trapping of triplet nitrene  $^3\mathbf{2}$ , as observed for the simple trifluoromethylnitrene [11] but not for methylnitrene [3, 4] remains unclear<sup>4</sup> [58, 59]. Addressing this question will likely require computing the corresponding reaction barriers on the singlet surface using multireference electronic structure methods.

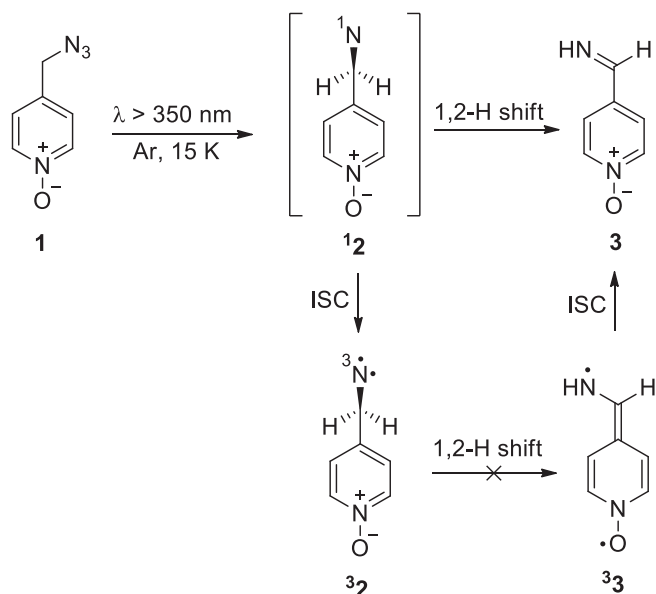
**TABLE 4** | Hirshfeld atomic spin population analysis of  $^3\mathbf{2}$  and  $^3\mathbf{3E}$  based on B3LYP/6-311+G(2d,p) calculations.

Atomic space <sup>a</sup>	Nitrene $^3\mathbf{2}$		Imine $^3\mathbf{3E}$	
	Value	(%)	Value	(%)
N1	1.67	83.7	0.36	18.0
C2	0.11	5.5	0.03	1.5
C3	0.05	2.7	0.32	15.9
C4	0.01	0.4	0.05	2.6
C5	0.00	0.0	0.14	7.2
N6	0.00	0.2	0.31	15.3
O7	0.02	0.8	0.58	28.8
C8	0.00	0.0	0.08	4.2
C9	0.01	0.4	0.08	4.2

<sup>a</sup>Atom numbering starts from the exocyclic nitrogen atom and follows the molecular connectivity in a clockwise direction.



**FIGURE 5** | Spin density plots of  $^3\mathbf{2}$  and  $^3\mathbf{3E}$  calculated at the B3LYP/6-311+G(2d,p) level of theory. Isosurfaces are shown at  $+0.01/-0.01 \text{ e bohr}^{-3}$  (orange/cyan). Color scheme: Carbon (gray), nitrogen (blue), oxygen (red), and hydrogen (white).



**SCHEME 2** | Summary of photolysis of 4-azidomethyl-pyridine-*N*-oxide **1** in an Ar matrix at 15 K. Proposed mechanistic pathways leading to the observed products, triplet nitrene **32** and imine **3** via common singlet nitrene **12** intermediate, in parallel with the mechanistic interpretation proposed by Platz et al. [11].

## 4 | Conclusions

This study provided unprecedented experimental access to the simplest triplet alkyl nitrene bearing a pyridine-*N*-oxide moiety, 4-nitrenemethyl-pyridine-*N*-oxide **32**, which was generated from photolysis of the 4-azidomethyl-pyridine-*N*-oxide **1** ( $\lambda > 350\text{ nm}$ ) isolated in an Ar matrix at 15 K. Alongside, 4-iminomethyl-pyridine-*N*-oxide **3** was also spectroscopically identified as a mixture of E and Z isomers. Monitoring the kinetics of the irradiation process shows that imine **3** does not arise from secondary photochemistry of triplet nitrene **32** but is likely formed via an initially unobserved transient singlet nitrene **12** in competition with the intersystem crossing to the triplet state.

Triplet nitrene **32** remained stable under cryogenic dark conditions, demonstrating that a putative quantum tunneling 1,2-hydrogen shift to the corresponding imine **3E** is not operative on the experimental timescale. Computational analysis reveals that the 1,2-H shift from **32** to **3E** is thermodynamically favored. The spin density in **32** is highly localized on the nitrene nitrogen atom, whereas in **33** is strongly delocalized, accounting for the unexpected relative stabilization of the latter. Nevertheless, this reaction involves a substantial energy barrier of  $\sim 36\text{ kcal mol}^{-1}$ , effectively precluding H-tunneling under cryogenic temperatures.

Overall, these results expand the fundamental understanding of alkyl nitrene chemistry and establish a well-defined model system for probing fundamental aspects of nitrene reactivity and tunneling phenomena. Furthermore, given that the pyridine-*N*-oxide moiety serves as an anchoring group capable of forming complexes with calix[4]pyrrole derivatives, these findings also open new perspectives for investigating alkyl nitrenes under supramolecular confinement and within host-guest architectures.

## Author Contributions

**Cláudio M. Nunes:** conceptualization, investigation, funding acquisition, writing – original draft, data curation, formal analysis. **A. J. Lopes Jesus:** formal analysis, writing – review and editing, data curation. **Tatiana Caneca:** investigation. **Gemma Aragay:** writing – review and editing, formal analysis, investigation. **Pablo Ballester:** writing – review and editing, formal analysis, funding acquisition. **Rui Fausto:** writing – review and editing, formal analysis, funding acquisition.

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## Conflicts of Interest

The authors declare no conflicts of interest.

## Data Availability Statement

The data that support the findings of this study are available in the [Supporting Information](#) of this article.

## Endnotes

<sup>1</sup> Some evidence of the triplet methyl nitrene, generated, for instance, via metastable energy-transfer fragmentation of methyl azide and methyl isocyanate in matrix-isolation experiments, has been reported based on EPR and UV-Vis spectroscopy; see Ferrante [5, 6].

<sup>2</sup> The generation of hydrogen-bonded complexes of triplet methyl nitrene in a cryogenic matrix has been reported recently. see Du et al. [7].

<sup>3</sup> Quantum tunneling probability decreases exponentially with increasing barrier width and with the square root of both the barrier

height and the mass of the tunneling particle. Consequently, an analysis based solely on energetic considerations cannot provide an accurate rationalization of tunneling reactivity. In the present case, the 1,2-H shift from triplet nitrene  $3^2$  to imine  $3^3$  is preceded by a significant reorganization of the heavy-atom framework, most notably a torsional motion of the nitrene fragment toward planarity, which occurs prior to hydrogen migration. Such heavy-atom rearrangements broaden the tunneling barrier and increase the effective mass along the reaction coordinate, thereby substantially reducing the tunneling probability at cryogenic temperatures. A quantitative prediction of tunneling probabilities would therefore require explicit multidimensional tunneling treatments, which are beyond the scope of this work.

<sup>4</sup>The pyrolysis of benzyl azide coupled with matrix isolation did not detect the corresponding triplet nitrene but only the imine derivatives. See Pinto et al. [58, 59].

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## Supporting Information

Additional supporting information can be found online in the Supporting Information section. **Figure S1:** Two-dimensional map of the potential energy of 4-azidomethyl-pyridine-*N*-oxide **1** as a function of the C–C–N and C–C–N–N dihedral angles. Each dihedral was incremented by 10°, with all remaining structural parameters fully optimized at the B3LYP/6-311+G(2d,p) level. Energy minima are denoted as **1g** (fourfold degenerate) and **1t** (doubly degenerate). The color scale corresponds to relative energies, referenced to the lowest minimum (**1g**), and isoenergy lines are drawn every 1 kJ mol<sup>-1</sup>. **Figure S2:** Potential energy diagram showing the relative zero-point-corrected energies of the **1g** and **1t** conformers of 4-azidomethyl-pyridine-*N*-oxide **1**, together with the transition state along the minimum energy path connecting the two conformers. The energy of the most stable conformer (**1g**) was taken as the zero of the relative energy scale. The energies shown were computed at the B3LYP/6-311+G(2d,p) (blue) and ωB97XD/6-311+G(2d,p) (red) levels of theory. **Figure S3:** Experimental IR difference spectra showing the spectral changes observed after 2 min of UV irradiation (λ ≥ 350 nm), followed by an additional 13 min of irradiation (15 min total) under identical conditions, of 4-azidomethyl-pyridine-*N*-oxide **1** isolated in an Ar matrix at 15 K. Positive and negative bands correspond to photogenerated and consumed species, respectively. The spectral changes indicate that the primary photoproducts **A** (assigned to triplet 4-nitrenemethyl-pyridine-*N*-oxide **3**) and **B** (assigned to 4-iminomethyl-pyridine-*N*-oxide **3**) are both formed within the first 2 min of irradiation. Upon prolonged irradiation, **B** is consumed concomitantly with **1**, whereas **A** continues to accumulate in the matrix. Bands marked with an asterisk are assigned unidentified secondary photoproducts. **Figure S4:** Time evolution of the amounts of 4-azidomethyl-pyridine-*N*-oxide **1** and of

its photoproducts 4-nitrenemethyl-pyridine-*N*-oxide **32** and the **E**- and **Z**-isomers of 4-iminomethyl-pyridine-*N*-oxide during UV exposure ( $\lambda \geq 350$  nm) of **1** isolated in an Ar matrix at 15 K for 15 min. The relative concentrations of reactant **1** and photoproducts **32**, **3E**, and **3Z** were estimated from the integrated areas of their characteristic infrared bands. Initially, diagnostic bands were selected based on their high intensity and absence of overlap with other spectral features:  $1346\text{ cm}^{-1}$  (**1**),  $1156\text{ cm}^{-1}$  (**32**),  $1143\text{ cm}^{-1}$  (**3E**), and  $1161\text{ cm}^{-1}$  (**3Z**). To account for the different transition dipole moments of these vibrations, the experimental integrated intensities were divided by the calculated intensities of the corresponding vibrations obtained from the B3LYP/6-311+G(2d,p) calculations (see Tables 2, 3, and S1). In the final step, these weighted integrated intensities were scaled by a common normalization factor, setting the initial amount of reactant **1** at  $t = 0$  min to 100%. At the early stages of photolysis ( $t \leq 2$  min), the sum of all identified species remains close to 100%. At longer irradiation times, this proportion decreases to 77% at  $t = 4$  min and to 62% at  $t = 15$  min, reflecting the formation of additional, unidentified secondary photoproducts, which account for the remaining 23% and 38%, respectively. **Figure S5**: Reaction energy profile ( $\Delta H_{0K}$  in  $\text{kcal mol}^{-1}$ ) computed at the  $\omega\text{B97XD}/6-311+\text{G}(2\text{d},\text{p})+\text{ZPVE}$  level of theory for the 1,2-H shift from triplet alkyl nitrenes to the corresponding imines. Values are relative to the energy of the corresponding triplet alkyl nitrenes. **Table S1**: Experimental IR spectral data for azide 4-azidomethyl-pyridine-*N*-oxide **1** isolated in an Ar matrix at 15 K, compared with the B3LYP/6-311+G(2d,p) vibrational frequencies ( $\nu$ ,  $\text{cm}^{-1}$ ) and infrared intensities ( $I$ ,  $\text{km mol}^{-1}$ ) calculated for the two conformers of **1** (**1g** and **1t**) and approximate vibrational assignments<sup>a</sup>.